

European Research Council



#### Nonthermal superconductivity

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#### Nonthermal superconductivity

in collaboration with:

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#### Motivation

"Tuning" of material properties by external driving

- Enhancement and control of electronic orders
  - e.g. light-induced high-temperature superconductivity (?)

cuprates: Fausti et al., Science (2011), Kaiser et al., PRB (2014)



phonons excited by THz pulse



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"Tuning" of material properties by external driving

- Enhancement and control of electronic orders
  - e.g. light-induced high-temperature superconductivity (?)

fullerides: Budden et al., Nature Phys. (2021)

stronger pump pulse (300fs - 300 ps)



#### Goal

- Explore nonthermal superconductivity in Hubbard models
  - I) Entropy cooling mechanism for producing "cold" photo-doped states
  - 2) eta-pairing in photo-doped Mott insulators on bipartite lattices
  - 3) Chiral superconductivity in photo-doped Mott insulators on frustrated lattices

- 5) Effective equilibrium approach for photo-doped Mott systems
- 6) Spin, charge and eta-spin separation in photo-doped ID systems

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numerical results based on dynamical mean field theory (DMFT)

5) Effective equilibrium approach for photo-doped Mott systems

6) Spin, charge and eta-spin separation in photo-doped ID systems

numerical results based on iterated time-evolving block decimation (iTEBD)



 Dynamical mean field theory DMFT: mapping to an impurity problem Georges & Kotliar, PRB (1992)

lattice model impurity model Δ DMFT self-consistency  $G_{\rm loc}^{\rm latt}(i\omega_n)$  $S_{\rm imp}[\Delta(i\omega_n)]$  $G_{\rm loc}^{\rm latt} \equiv G_{\rm imp}$ momentum average impurity solver  $\Sigma_k^{\text{latt}} \equiv \Sigma_{\text{imp}}$  $\frac{1}{i\omega_n + \mu - \epsilon_k - \Sigma_{\nu}^{\text{latt}}}$  $G_k^{\text{latt}}$  $G_{\rm imp}(i\omega_n), \Sigma_{\rm imp}(i\omega_n)$ 

DMFT approximation

#### Conceptual question

Appearance of "fragile" electronic orders in highly nonequilibrium states



correlated electron system

### Conceptual question

Appearance of "fragile" electronic orders in highly nonequilibrium states

Need to find ways to avoid heating

entropy cooling



correlated electron system

- Photo-doping from/to flat bands Werner, Eckstein, Mueller & Refael, Nat. Comm. (2019)
  - Dipolar excitations with appropriate frequency  $\Omega$  transfer electrons from core to system and cool down the system



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- Photo-doping from/to flat bands Werner, Eckstein, Mueller & Refael, Nat. Comm. (2019)
  - Entropy of the core band in the narrow band (atomic) limit:

$$S_{\text{core}} = -2n_{\sigma}\ln(n_{\sigma}) - 2(1-n_{\sigma})\ln(1-n_{\sigma})$$



In case of isentropic doping process:

 $\Delta S_{\rm core} \nearrow \Rightarrow \Delta S_{\rm system} \searrow$ 

cooling of system due to entropy reshuffling

- Photo-doping from/to flat bands Werner, Eckstein, Mueller & Refael, Nat. Comm. (2019)
  - Constant entropy contours in the filling-temperature plane



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  - Constant entropy contours in the filling-temperature plane



AFM order

## Entropy trapping

- Thermalization bottleneck prevents system from heating
  - e.g. recombination of electrons and holes can be very slow if the gap size is large



Nonthermal superconductivity in entropy-cooled systems

•  $\eta$  pairing in a repulsive Hubbard model with inverted population and positive effective doublon/holon temperature

Rosch, Rasch, Binz & Vojta, PRL (2008); Werner, Li, Golez & Eckstein, PRB (2019)



state with almost complete population inversion and "cold" effective T>0 prepared by entropy-cooling protocol

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Nonthermal superconductivity in entropy-cooled systems

- Nonequilibrium phase diagram of photo-doped Mott insulators
- Use steady-state DMFT to control doublon concentration and  $T_{eff}$

Li, Golez, Werner & Eckstein, PRB (2020)



steady-state formalism

- Nonthermal superconductivity in entropy-cooled systems
  - Nonequilibrium phase diagram of photo-doped Mott insulators
- (a) • Use steady-state DMFT to control doublon concentration and  $T_{\rm eff}$ Li, Golez, Wegner & Eckstein, PRB (2020) full fermion bath  $10^{4}$ 20 ¦afm n-pairing upper Hubbard 15  $\chi_\eta$ •  $10^{3}$ 10 doublons 5 10<sup>2</sup> 6 -6 -4 -2 0 2 4 8 recombination  $\beta_{eff}$ 0 0.4  $10^{1}$ (b) ъSC -5 holes n order -10 0.3 empty fermion bath appo 10<sup>0</sup> ferromagnetism s-wave SC -15 CDW lower Hubbard  $10^{-1}$ -20 0.1 0.2 0.3 0.4 0.5 0 d 0.1 steady-state formalism phase diagram 0.0

Nonthermal superconductivity in entropy-cooled systems

 $\bullet$  Optical conductivity of the  $\eta$  pairing state

Li, Golez, Werner & Eckstein, PRB (2020)



- No gap in the real part (in contrast to s-wave SC)
- two-fluid picture: condensed doublons/holons coexist with normal singlons

Nonthermal superconductivity in entropy-cooled systems

- What happens in photo-doped systems on frustrated lattices?
- Can we realize an analogue of 120-degree order?

Li, Mueller, Kim, Laeuchli & Werner, arXiv (2022)



chiral superconducting state with loop currents

experimental signature:

second-order transverse supercurrent response for A along x

Nonthermal superconductivity in entropy-cooled systems

- What happens in photo-doped systems on frustrated lattices?
- Can we realize an analogue of I20-degree order?

Li, Mueller, Kim, Laeuchli & Werner, arXiv (2022)

 $\bullet\,$  Evidence from entropy-cooling protocol (for  $\varphi=0)$ 



Nonthermal superconductivity in entropy-cooled systems

- What happens in photo-doped systems on frustrated lattices?
- Can we realize an analogue of 120-degree order?

Li, Mueller, Kim, Laeuchli & Werner, arXiv (2022)

• Evidence from exact diagonalization (12 sites)



Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

• U-V Hubbard model

$$H = -t_{\rm hop} \sum_{\langle i,j\rangle,\sigma} (c_{i\sigma}^{\dagger}c_{j\sigma} + h.c.) + H_U + H_V$$

• Photo-doping leads to steady-state with "cold" doublons/holons



#### Compute effective Hamiltonian by Schrieffer-Wolff transformation

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$$H = -t_{\rm hop} \sum_{\langle i,j\rangle,\sigma} (c_{i\sigma}^{\dagger}c_{j\sigma} + h.c.) + H_U + H_V$$

• Eliminate terms which change number of doublons/holons/singlons

$$H_{\text{eff}} = H_U + H_{\text{kin,doublon}} + H_{\text{kin,holon}} + H_V + H_{\text{spin-ex}} + H_{\text{dh-ex}} + H_{\text{U-shift}} + H_{3\text{-site}}$$

$$H_{\text{spin-ex}} = J_{\text{ex}} \sum_{\langle i,j \rangle} \vec{s}_i \cdot \vec{s}_j$$

spin exchange term determines correlations between neighobring singlons

$$H_{\text{dh-ex}} = -J_{\text{ex}} \sum_{\langle i,j \rangle} \vec{\eta}_i \cdot \vec{\eta}_j$$

doublon-holon exchange term determines correlations between neighboring doublon-holon pairs

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$$H_{\text{eff}} = H_U + H_{\text{kin,doublon}} + H_{\text{kin,holon}} + H_V + H_{\text{spin-ex}} + H_{\text{dh-ex}} + H_{\text{U-shift}} + H_{3\text{-site}}$$

$$H_{\text{spin-ex}} = J_{\text{ex}} \sum_{\langle i,j \rangle} \vec{s}_i \cdot \vec{s}_j \qquad \eta_i^+ = (-1)^i c_{i\downarrow}^\dagger c_{i\uparrow}^\dagger \eta_i^- = (-1)^i c_{i\uparrow} c_{i\downarrow} \eta_i^- = (-1)^i c_{i\uparrow} c_{i\downarrow} \eta_i^z = \frac{1}{2} (n_i - 1)$$

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$$H = -t_{\rm hop} \sum_{\langle i,j\rangle,\sigma} (c_{i\sigma}^{\dagger}c_{j\sigma} + h.c.) + H_U + H_V$$

• Eliminate terms which change number of doublons/holons/singlons

$$\begin{split} J_{\mathrm{ex}} &= \frac{4t_{\mathrm{hop}}^2}{U} \\ (\uparrow,\downarrow) &\to \underbrace{(\uparrow\downarrow,0)}_{\Delta E=U} \to (\downarrow\uparrow) \qquad \text{spin exchange is antiferro} \\ (\uparrow\downarrow,0) &\to \underbrace{(\uparrow,\downarrow)}_{\Delta E=-U} \to (0,\uparrow\downarrow) \qquad \text{doublon-holon exchange is ferro} \end{split}$$

Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

Introduce separate chemical potentials for doublons and holons

$$N_{h} = \sum_{i} n_{i}^{h}, \quad n_{i}^{h} = (1 - n_{i\uparrow})(1 - n_{i\downarrow})$$
$$N_{d} = \sum_{i} n_{i}^{d}, \quad n_{i}^{d} = n_{i\uparrow}n_{i\downarrow}$$

Introduce grand-canonical Hamiltonian for photo-doped state

Then use favorite equilibrium method to solve this problem

Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

- Static observables can be computed directly from  $ho_{
  m eff} = e^{-eta_{
  m eff}K_{
  m eff}}$
- Response functions  $-i\langle [A(t), B(0)]_{\pm} \rangle$  can also be computed, but
  - ullet initial state described by  $K_{\mathrm{eff}}$
  - ullet time propagation determined by  $H_{
    m eff}$

must split the operators A and B into terms which change the doublon number by +1, 0, -1 and multiply these terms with appropriate phase factors

$$A = \sum_{\alpha} A_{\alpha} \quad H_{\mu} = -\sum_{G} \mu_{G} N_{G}$$
$$\lambda_{\alpha} = -\sum_{G} \mu_{G} \Delta N_{G,\alpha}$$
$$e^{-iH_{\mu}t} A_{\alpha} e^{iH_{\mu}t} = e^{-i\lambda_{\alpha}t} A_{\alpha}$$

 To determine the nonequilibrium phase diagram, measure the decay of spin, charge and eta-spin correlations, e.g. using iTEBD

Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

• Charge and eta-pairing correlations in the photo-doped ID U-V Hubbard model ( $H_{eff}$  without 3-site terms) at  $n_d = n_h = 0.23$ 



charge correlations dominate for V=0.4

eta pairing correlations dominate for V=0.1, 0.2

Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

• "Zero effective temperature" phase diagram of photo-doped ID U-V Hubbard model ( $H_{eff}$  without 3-site terms)



Compute effective Hamiltonian by Schrieffer-Wolff transformation

Murakami et al., Comm. Physics (2022)

• Spectral functions of the photo-doped ID U-V Hubbard model  $(H_{\rm eff}$  without 3-site terms)



not gapped for the eta-pairing phase

- Exact wave function in the limit of large on-site repulsion Murakami et al., arxiv:2212.06263 (2022)
  - Consider  $H_{\rm eff}$  with fixed number of doublons and holons
  - Wave function in the limit  $J_{ex} \rightarrow 0, V/J_{ex} = const$  is a generalization of the Ogata-Shiba state

Ogata & Shiba, PRB 41, 2326 (1990) (doped equilibrium model)

• For  $J_{\rm ex} = V = 0$ : Eigenstates of  $H_{\rm eff}$  are degenerate w.r.t. spin and eta-spin configurations

 $H_{\rm kin, doublon} + H_{\rm kin, holon}$  does not flip spins or exchange d-h pairs

$$|\Psi\rangle = |\Psi_{\rm SF}^{\rm GS}\rangle|\Psi_{\sigma,\eta}\rangle$$

ground state of spinless Fermions

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$$|\Psi\rangle = |\Psi_{\rm SF}^{\rm GS}\rangle |\Psi_{\sigma,\eta}\rangle \leftarrow \text{degeneracy of } 2^{N_s} 2^{N_\eta} \text{ lifted by } \mathcal{O}(J_{\rm ex}) \text{ terms}$$

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• After taking into account the  $\mathcal{O}(J_{ex})$  terms, the squeezed spin and eta-spin spaces get decoupled

$$H_{\text{spin}}^{\text{squeezed}} = J_{\text{ex}}^s \sum_i \vec{s}_{i+1} \cdot \vec{s}_i$$

$$H_{\eta-\text{spin}}^{\text{squeezed}} = -J_X^s \sum_j (\eta_{j+1}^x \eta_j^x + \eta_{j+1}^y \eta_j^y) + J_Z^s \sum_j \eta_{j+1}^z \eta_j^z$$

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 $|\Psi\rangle = |\Psi_{\rm SF}^{\rm GS}\rangle |\Psi_{\sigma}^{\rm GS}\rangle |\Psi_{\eta}^{\rm GS}\rangle$ 

eta-pairing phase has C=3

CDW phase has C=2

confirmed by iTEBD analysis of the entanglement entropy



#### Conclusions

- Nonthermal superconducting states in the Hubbard model
  - Mechanisms for realizing cold photo-doped Mott states
    - Entropy cooling: photo-doping from flat bands
    - Coupling to baths: injection of "cold" doublons and holes
  - Examples
    - eta-pairing on bipartite lattices
    - chiral superconductivity on geometrically frustrated lattices
  - Spin, charge and eta-spin separation in photo-doped ID Mott systems

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