# Towards a mathematical theory of the ODE／IM correspondence 

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# Integrability in Condensed Matter and Quantum Field Theories 

SWISS MAP Les Diablerets， 7 February 2023

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## References

The talk is based on three recent papers with R. Conti and $A$. Raimondo:

- R. Conti and D.M., On solutions of the Bethe Ansatz for the Quantum KdV model. arXiv 2022
- R. Conti and D.M., Counting Monster Potentials. JHEP 2021
- D.M. and Andrea Raimondo, Opers for higher states of quantum KdV models, Commm. Math. Phys, 2020.
And ongoing work with G. Degano, E. Mukhin, and A. Raimondo.


## Plan of the talk

- Introduction
- If the momentum is large enough, the Destri-De Vega equation for Quantum KdV is well-posed.
- The monster potentials are complete (proven up to a technical assumption)


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## Linear ODEs and Math Phys

It is a general fact that many interesting objects in mathematical physics can be expressed via the (generalised) monodromy data of a of linear ODE in the complex plane.

Quantum field theories \& related geometrical objects often
correspond to equations with irregular singularities (e.g. Frobenius manifolds and TQFT, ODE/IM correspondence).
This correspondence is somehow mediated by the Nonlinear Stokes Phenomenon (Wall-Crossing).

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## (About names)

'ODE/IM correspondence' and 'ODE/IQFT correspondence' are two different names for the same thing.

We should all agree to use the same name.

## Bethe Equations for Quantum KdV

Fix $\alpha>0, p \geq 0$ (in most of this talk, $\alpha>1$ ).
We look for a real entire function $Q(E)$ of order $\frac{1+\alpha}{2 \alpha}$ such that

- All roots are simple, zero is not a root $(Q(0)=1)$.
- (Almost) all roots are real and positive.
- The counting function $n(E)$ satisfies $\lim _{E \rightarrow+\infty} \frac{n(E)}{E^{\frac{1+\alpha}{2 \alpha}}}=1$.
- If $E_{*}$ is a root, then

$$
e^{-i 4 \pi p} \frac{Q\left(e^{-\frac{2 \pi i}{\alpha+1}} E_{*}\right)}{Q\left(e^{\frac{2 \pi i}{\alpha+1}} E_{*}\right)}=-1
$$

Additional hypothesis: number of holes are finite.

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Additional hypothesis: number of holes are finite.

$$
Q\left(E_{*}\right)=0 \Longrightarrow e^{-i 4 \pi p} \frac{Q\left(e^{-\frac{2 \pi i}{\alpha+1}} E_{*}\right)}{Q\left(e^{\frac{2 \pi i}{\alpha+1}} E_{*}\right)}=-1
$$

- Edge asymptotics of Bethe roots for XXZ with $-1<\Delta<1$ (cf. H. Boos talk).
- Bethe Equations satisfied by the eigenvalues of the $\mathcal{Q}_{+}$ operator of Quantum KdV
(Bazhanov-Lukyanov-Zamolodchikov '94-'96).

$$
c=1-\frac{6 \alpha^{2}}{\alpha+1}, \Delta=p^{2}(1+\alpha)-\frac{\alpha^{2}}{4 \alpha+4} .
$$

- The same equations are also satisfied by spectral determinants of some anharmonic oscillator (Dorey-Tateo '98 \& BLZ '98).

- Problem: Classify solutions of the BE whose roots are all real and positive.
- Introducing the 'counting function',

$$
z(E)=-2 p+\frac{1}{2 \pi i} \log \frac{Q\left(e^{-i \frac{2 \pi}{\alpha+1}} E\right)}{Q\left(e^{i \frac{2 \pi}{\alpha+1}} E\right)}, E \geq 0, z(0)=-2 p
$$

- BE reads $Q\left(E_{*}\right)=0 \Longrightarrow z\left(E_{*}\right)-\frac{1}{2} \in \mathbb{Z}$
- To classify solution we must add information on which quantum numbers are occupied.


## Roots and Holes

- $E_{k}: z\left(E_{k}\right)=k+\frac{1}{2}$ with $k \in \mathbb{Z}, k \geq-2 p+\frac{1}{2}$
- $k$ is a root number if $Q\left(E_{k}\right)=0$, a hole-number otherwise.
- Root numbers form as increasing sequence $\left\{k_{n}\right\}_{n \in \mathbb{N}}$
- Now the BE reads

$$
z\left(E_{k_{n}}\right)=k_{n}+\frac{1}{2}, n \in \mathbb{N}, \text { with } Q(E)=\prod_{n}\left(1-\frac{E}{E_{k_{n}}}\right) .
$$

We want to study well-posedness of the above equation, when the sequence $\left\{k_{n}\right\}_{n \in \mathbb{N}}$ is given.

## Roots and integer partitions

Root-numbers are sequences that stabilizes: $k_{n}=n$, if $n \gg 0$.
$\Downarrow$
Root-numbers sequences are classified by integer partitions $\left\{k_{n}^{\lambda}\right\}_{n \in \mathbb{N}}$.


## The ODE/IM Conjecture for Quantum KdV

Bazhanov-Lukyanov-Zamolodchikov, Adv. Theor. Math. Phys., (2003) made the following conjecture:
(1) Let $N \in \mathbb{N}$ and $2 p \geq N+\frac{1}{2}$. For every $\lambda \vdash N$, the BE admit a unique (normalised) solution $Q^{\lambda}(E ; p)$ whose sequence of root-numbers coincide with $\left\{k_{n}^{\lambda}\right\}_{n \in \mathbb{N}}$.
(3) Any such solution of the BE coincides with the spectral determinant of a certain anharmonic oscillator.

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(2) Any such solution of the BE coincides with the spectral determinant of a certain anharmonic oscillator.

## Our results. 1. Well-posedness of BE

## (1) Theorem, M. - Conti 2022

Fix $\alpha>1,(N, \lambda \vdash N)$.
If $p$ is sufficiently large:

1. The BE admit a unique solution $Q^{\lambda}(E ; p)$ whose sequence of root-numbers coincide with $\left\{k_{n}^{\lambda}\right\}_{n \in \mathbb{N}}$.
2. $\forall n \in \mathbb{N}, \exists C_{n}>0$ such that

$$
\left|\frac{E_{k_{n}}(p)}{p^{\frac{2 \alpha}{1+\alpha}}}-\left[A+B\left(k_{n}+\frac{1}{2}\right) \frac{1}{p}\right]\right| \leq \frac{C_{n}}{p^{2}} .
$$

3. Uniform asymptotics of $z$ and of roots.

## Earlier results in the mathematical literature

- Well-posedness for $\alpha>1, p=\frac{1}{2 \alpha+2}$ and $\lambda=\emptyset$ by A. Avila in Comm. Math. Phys. (2004) - after Voros.
- Well-posedness for $2 \alpha$ integer and $\lambda=\emptyset$ by Hilfiker and Runke, Ann. Henri Poincaré (2020), using TBA.

> Remark. A variational approach (à la Yang \& Yang) should yield sharp bound on the range of $p$ for which BE with real roots only is well-posed.

The range $0<\alpha<1$ seems more difficult to study.

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## Strategy of the proof

- Transform the logarithmic BE into a Free-Boundary Nonlinear Integral Equation (known as Destri-De Vega).
- Linearise in the large $p$ limit and do perturbation analysis.

The strategy is standard, the analysis is completely new.

## Destri-De Vega Integral Equation

Given $\lambda \vdash N$, let $H=\#\{$ holes greater than the lowest root $\}$. The unknown is a tuple $\left(\omega, h_{1}, \ldots, h_{H}, z\right)$

- $\omega>0$, the left endpoint of the integration interval $[\omega,+\infty[$;
- $h_{1}<\cdots<h_{H}$ are the holes greater than the lowest root;
- $z$ : $C^{1}\left(\left[\omega, \infty[)\right.\right.$, strictly monotone, $z(E) \sim E^{\frac{1+\alpha}{2 \alpha}}, x \rightarrow+\infty$.


## Destri-De Vega Integral Equation II

The Destri-De Vega (DDV) is a free-boundary nonlinear integral equation:

$$
\begin{aligned}
& \text { 1. } z(E)=-2 p+\int_{\omega}^{\infty} K_{\alpha}(E / y)\left\lceil z(y)-\frac{1}{2}\right\rceil \frac{d y}{y}+H F_{\alpha}\left(\frac{E}{\omega}\right) \\
& -\sum_{j=1}^{H} F_{\alpha}\left(\frac{E}{h_{k}}\right), K_{\alpha}(x)=x F_{\alpha}^{\prime}(x)=\frac{\sin \left(\frac{2 \pi}{1+\alpha}\right)}{\frac{1}{\pi}} \frac{x}{1+x^{2}-2 x \cos \left(\frac{2 \pi}{1+\alpha}\right)} \\
& \text { 2. }\left\lceil z(\omega)-\frac{1}{2}\right\rceil=-H \\
& \text { 3. } z\left(h_{j}\right)=\sigma(j)+\frac{1}{2}, j=1 \ldots H, \sigma(j)=\text { quantum number of } h_{j} \\
& \text { Remark. If } z \text { is a strictly increasing real analytic function }
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Remark. If $z$ is a strictly increasing real analytic function

$$
\lim _{\varepsilon \rightarrow 0^{+}} \frac{1}{\pi} \operatorname{lm} \log \left(1+e^{2 \pi i z(x+i \varepsilon)}\right)=z-\left\lceil z-\frac{1}{2}\right\rceil
$$

## Large $p$ Linearisation $=$ WKB

$$
z_{\omega, p}(E)=-2 p+\int_{\omega}^{\infty} K_{\alpha}(E / y) z_{\omega, p}(y) \frac{d y}{y}, z_{\omega, p}(E) \sim E^{\frac{\alpha+1}{2 \alpha}}, x \rightarrow \infty
$$

It is a Wiener-Hopf equation, solutions can be expressed via

$$
\tau(\xi)=\frac{1}{2 \pi i} \int_{\delta-i \infty}^{\delta+i \infty} \frac{\frac{\alpha s}{1+\alpha}}{2 \sqrt{\pi}(1+\alpha)^{s-1}}\left\ulcorner\frac{\Gamma\left(-\frac{1}{2}-\frac{\alpha s}{1+\alpha}\right) \Gamma\left(1-\frac{s}{1+\alpha}\right)}{s^{2}\ulcorner(-s)} \xi^{-s} d s, \quad \xi=x / \omega .\right.
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We discovered a (much more useful) formula in terms of a WKB integral


This is a first hint of the ODE/IM correspondence.

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$$
\tau(\xi)=\frac{1}{\pi} \int_{u_{-}}^{u_{+}} \sqrt{u^{2} \xi-u^{2 \alpha+2}-1} \frac{d u}{u}, \sqrt{\cdots} \mid u=u_{ \pm}=0 .
$$

This is a first hint of the ODE/IM correspondence.

We need to analyse integrals like

$$
\begin{aligned}
& A_{p}[f, \varepsilon]=\int_{1}^{\infty} K_{\alpha}\left(\frac{x}{y}\right)\langle p f(y)+\varepsilon(y)\rangle \frac{d y}{y},\langle z\rangle=z-\left\lceil z-\frac{1}{2}\right\rceil \\
& B_{p}[f, \varepsilon]=\int_{1}^{\infty} K_{\alpha}\left(\frac{x}{y}\right)\left\lceil p f(y)+\varepsilon(y)-\frac{1}{2}\right\rceil \frac{d y}{y}
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As an example, we showed that if $f \sim x^{\frac{\alpha+1}{2 \alpha}}$ and $\varepsilon, \tilde{\varepsilon}$ are bounded ( + some further hypotheses), then


## $\Longrightarrow$ contractiveness of the perturbation operator $B_{p}[I, \cdot]$ when $p$ is

 large.
## Perturbation/Analytical challenges

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$$
\left|\left\|B_{p}[f, \varepsilon]-B_{p}[f, \tilde{\varepsilon}]\right\|_{\infty}-\frac{\alpha+1}{2 \alpha}\|\varepsilon-\tilde{\varepsilon}\|_{\infty}\right| \lesssim f \frac{\|\varepsilon-\tilde{\varepsilon}\|_{\infty}}{p}
$$

$\Longrightarrow$ contractiveness of the perturbation operator $B_{p}[/, \cdot]$ when $p$ is large.

$$
-\Psi^{\prime \prime}(x)+\left(x^{2 \alpha}+\frac{\ell(\ell+1)}{x^{2}}-E\right) \Psi(x)=0, \alpha>1, \ell \geq 0, E \in \mathbb{C}
$$

$E$ is said an eigenvalue if $\exists \Psi \neq 0$ such that

$$
\lim _{x \rightarrow 0^{+}} \Psi(x)=\lim _{x \rightarrow+\infty} \Psi(x)=0
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The spectrum is discrete, simple and positive, $E_{n}(\ell), n \in \mathbb{N}$ :


Spectral determinant $D_{\ell}(E)$ is an entire function of order $\frac{1+\alpha}{2 \alpha}$

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$$
E_{n}(\ell)=\left(\frac{2 \Gamma\left(\frac{2 \alpha+1}{2 \alpha}\right)}{\sqrt{\pi} \Gamma\left(\frac{3 \alpha+1}{2 \alpha}\right)}\right)^{-\frac{2 \alpha}{\alpha+1}}(4 n+2 \ell+3)^{\frac{2 \alpha}{\alpha+1}}\left(1+O\left(n^{-1}\right)\right)
$$

Spectral determinant $D_{\ell}(E)$ is an entire function of order $\frac{1+\alpha}{2 \alpha}$.

## Monster potentials, BLZ (2003)

1. Let $R$ be a monic polynomial of degree $N$. The spectral determinant $D_{\ell}^{R}(E)$ for the potential

$$
V^{R}=x^{2 \alpha}+\frac{\ell(\ell+1)}{x^{2}}-2 \frac{d^{2}}{d x^{2}} \log R\left(x^{2 \alpha+2}\right)
$$

satisfies the BE if the monodromy about the additional poles is trivial for every $E$.
2. Assuming that the roots of $R$ are distinct, the trivial monodromy is equivalent to the BLZ system

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2. Assuming that the roots of $R$ are distinct, the trivial monodromy is equivalent to the BLZ system

$$
\sum_{j \neq k} \frac{z_{k}\left(z_{k}^{2}+(3+\alpha)(1+2 \alpha) z_{k} z_{j}+\alpha(1+2 \alpha) z_{j}^{2}\right)}{\left(z_{k}-z_{j}\right)^{3}}-\frac{\alpha z_{k}}{4(1+\alpha)}+\Delta(\ell, \alpha)=0, \quad k=1, \ldots, N .
$$

## Wronskian of Hermite polynomials

## Rational extensions of the harmonic oscillator

- A rational extension of degree $N$ is a potential

$$
V^{U}(t)=t^{2}-2 \frac{d^{2}}{d t^{2}} \ln U(t)
$$

where $U$ a polynomial of degree $N$ such that all monodromies of $\psi^{\prime \prime}(t)=\left(V^{U}(t)-E\right) \psi$ are trivial for every $E$.

- Oblomkov's theorem (1999)

$$
U \propto U^{\lambda}:=W r\left[H_{\lambda_{1}+j-1}, \ldots, H_{\lambda_{j}}\right], \text { for a } \lambda:=\left(\lambda_{1}, \ldots, \lambda_{j}\right) \vdash N .
$$

## Large momentum limit of Monster Potentials

## (2) (Conditional) Theorem, M. - Conti 2021/2022

- We noticed that in the large momentum multi-scale limit, monster potentials converge to rational extensions of the harmonic oscillator:
Assume there exists a sequence $R_{\ell}$ of monster potentials with $\ell \rightarrow \infty$, then - up to subsequences -

$$
z_{k}=\frac{\ell^{2}}{\alpha}+\frac{(2 \alpha+2)^{\frac{3}{4}}}{\alpha} v_{k}^{\lambda} \ell^{\frac{3}{2}}+O(\ell), k=1, \ldots, N
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where $v_{k}^{\lambda}$ are the roots of $U^{\lambda}$.


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$$

where $v_{k}^{\lambda}$ are the roots of $U^{\lambda}$.

- (If a monster potential with a such an asymptotics exists and) $D_{\ell}^{\lambda}(E)$ is the corresponding spectral determinant, then

$$
D^{\lambda}(E ; \ell)=Q^{\lambda}(E / \eta ; p), p=\frac{2 \ell+1}{\alpha+1} \text { and } \eta=\left(\frac{2 \sqrt{\pi} r\left(\frac{3}{2}+\frac{1}{2 \alpha}\right)}{\Gamma\left(1+\frac{1}{2 \alpha}\right)}\right)^{\frac{2 \alpha}{1+\alpha}}
$$

## An unproven identity

Let $\lambda \vdash N$, assume $U^{\lambda}$ has $N$ distinct zeroes (see conjecture by Felder-Hemery-Veselov 2010). Consider the Jacobian

$$
J_{i j}^{\lambda}=\delta_{i j}\left(1+\sum_{l \neq j} \frac{6}{\left(v_{i}^{\lambda}-v_{j}^{\lambda}\right)^{4}}\right)-\left(1-\delta_{i j}\right) \frac{6}{\left(v_{i}^{\lambda}-v_{j}^{\lambda}\right)^{4}}, i, j=1, \ldots, N .
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$$

We found an explicit (albeit unproven) formula for the eigenvalues of $J^{\lambda}$ : these are square of the hook-lengths of the partition.
$\lambda=(N)$ stated/proven in Ahmed, Bruschi, Calogero, Olshanetsky, and Perelomov ('79).

## The Big ODE/IM Conjecture

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If a QFT is Bethe Integrable then the corresponding solutions of the Bethe Equations are spectral determinants of linear differential operators.
$\rightarrow$ Bethe Roots are eigenvalues of a (possibly self-adjoint) differential operator (cf. Hilbert-Pólya Conjecture).


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M - Raimondo (- Valeri) ('16,'17, '20, ongoing) after
Feigin-Frenkel (2011)
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$\widehat{\mathfrak{g}}$ an affine Kac-Moody Lie-algebra and ${ }^{L} \widehat{\mathfrak{g}}$ the Langlands dual,
$\{$ Bethe states of $\widehat{\mathfrak{g}}$ - quantum KdV$\} \leftrightarrow \cdots \rightarrow\left\{\widehat{\mathfrak{g}}-\right.$ opers on $\left.\mathbb{C}^{*}\right\}$.

## Open questions? A lot

The ODE/IM correspondence for Quantum KdV is just a tiny piece of an enormous field of research of which we know a lot but still very little.
> - How do we guess which ODE (if any) corresponds to a given Quantum Field Theory?
> - Once, they are found, how do we prove them?
> - Why the ODE/IM correspondence? Can we find a theory? Why is the nonlinear Stokes phenomenon that important?

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