Towards a mathematical theory of the ODE/IM correspondence

Davide Masoero Grupo de Física Matemática da Universidade de Lisboa With the support of the FCT Project UIDB/00208/2020

Integrability in Condensed Matter and Quantum Field Theories

SWISS MAP Les Diablerets, 7 February 2023





The talk is based on three recent papers with R. Conti and A. Raimondo:

- R. Conti and D.M., *On solutions of the Bethe Ansatz for the Quantum KdV model.* arXiv 2022
- R. Conti and D.M., Counting Monster Potentials. JHEP 2021
- D.M. and Andrea Raimondo, *Opers for higher states of quantum KdV models*, Commm. Math. Phys, 2020.

And ongoing work with G. Degano, E. Mukhin, and A. Raimondo.

Introduction

• If the momentum is large enough, the Destri-De Vega equation for Quantum KdV is well-posed.

• The monster potentials are complete (proven up to a technical assumption).

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It is a general fact that many interesting objects in mathematical physics can be expressed via the (generalised) monodromy data of a of linear ODE in the complex plane.

Quantum field theories & related geometrical objects often correspond to equations with irregular singularities (e.g. Frobenius manifolds and TQFT, ODE/IM correspondence). This correspondence is somehow mediated by the Nonlinear Stokes Phenomenon (Wall-Crossing).

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'ODE/IM correspondence' and 'ODE/IQFT correspondence' are two different names for the same thing.

We should all agree to use the same name.

Fix $\alpha > 0, p \ge 0$ (in most of this talk, $\alpha > 1$). We look for a real entire function Q(E) of order $\frac{1+\alpha}{2\alpha}$ such that

- All roots are simple, zero is not a root (Q(0) = 1).
- (Almost) all roots are real and positive.
- The counting function n(E) satisfies $\lim_{E \to +\infty} \frac{n(E)}{E^{\frac{1+\alpha}{2\alpha}}} = 1$.
- If E_{*} is a root, then

$$e^{-i4\pi p}\frac{Q\left(e^{-\frac{2\pi i}{\alpha+1}}E_*\right)}{Q\left(e^{\frac{2\pi i}{\alpha+1}}E_*\right)}=-1.$$

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Additional hypothesis: number of holes are finite.

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Where do they appear?

$$Q(E_*) = 0 \Longrightarrow e^{-i4\pi p} rac{Qig(e^{-rac{2\pi i}{lpha+1}}E_*ig)}{Qig(e^{rac{2\pi i}{lpha+1}}E_*ig)} = -1.$$

- Edge asymptotics of Bethe roots for XXZ with $-1 < \Delta < 1$ (cf. H. Boos talk).
- Bethe Equations satisfied by the eigenvalues of the Q₊ operator of Quantum KdV (Bazhanov-Lukyanov-Zamolodchikov '94-'96).

$$c=1-rac{6lpha^2}{lpha+1}, \Delta=
ho^2(1+lpha)-rac{lpha^2}{4lpha+4}.$$

 The same equations are also satisfied by spectral determinants of some anharmonic oscillator (Dorey-Tateo '98 & BLZ '98).

The ODE/IM Conjecture for Quantum KdV



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Topological classification of purely real solutions

- Problem: <u>Classify solutions of the BE</u> whose roots are **all** real and positive.
- Introducing the 'counting function',

$$z(E) = -2p + rac{1}{2\pi i}\lograc{Qig(e^{-irac{2\pi}{lpha+1}}Eig)}{Qig(e^{irac{2\pi}{lpha+1}}Eig)}, E\geq 0, z(0) = -2p.$$

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- BE reads $Q(E_*) = 0 \Longrightarrow z(E_*) \frac{1}{2} \in \mathbb{Z}$
- To classify solution we must add information on which quantum numbers are occupied.

Roots and Holes

•
$$E_k: z(E_k) = k + \frac{1}{2}$$
 with $k \in \mathbb{Z}, k \ge -2p + \frac{1}{2}$

- k is a root number if $Q(E_k) = 0$, a hole-number otherwise.
- Root numbers form as increasing sequence $\{k_n\}_{n\in\mathbb{N}}$
- Now the BE reads

$$z(E_{k_n})=k_n+rac{1}{2},\ n\in\mathbb{N},\ ext{with}\ Q(E)=\prod_n\left(1-rac{E}{E_{k_n}}
ight).$$

We want to study well-posedness of the above equation, when the sequence $\{k_n\}_{n \in \mathbb{N}}$ is given.

Roots and integer partitions

Root-numbers are sequences that stabilizes: $k_n = n$, if $n \gg 0$.

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Root-numbers sequences are classified by integer partitions $\{k_n^{\lambda}\}_{n \in \mathbb{N}}$.



Bazhanov-Lukyanov-Zamolodchikov, Adv. Theor. Math. Phys., (2003) made the following conjecture:

Let N ∈ N and 2p ≥ N + ¹/₂. For every λ ⊢ N, the BE admit a unique (normalised) solution Q^λ(E; p) whose sequence of root-numbers coincide with {k^λ_n}_{n∈N}.

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(1) Theorem, M. - Conti 2022

Fix $\alpha > 1$, $(N, \lambda \vdash N)$.

If p is sufficiently large:

1. The BE admit a unique solution $Q^{\lambda}(E; p)$ whose sequence of root-numbers coincide with $\{k_n^{\lambda}\}_{n \in \mathbb{N}}$.

2. $\forall n \in \mathbb{N}, \exists C_n > 0$ such that

$$\left|\frac{E_{k_n}(p)}{p^{\frac{2\alpha}{1+\alpha}}}-\left[A+B\left(k_n+\frac{1}{2}\right)\frac{1}{p}\right]\right|\leq \frac{C_n}{p^2}.$$

3. Uniform asymptotics of z and of roots.

Earlier results in the mathematical literature

- Well-posedness for $\alpha > 1$, $p = \frac{1}{2\alpha+2}$ and $\lambda = \emptyset$ by A. Avila in Comm. Math. Phys. (2004) after Voros.
- Well-posedness for 2α integer and λ = Ø by Hilfiker and Runke, Ann. Henri Poincaré (2020), using TBA.

Remark. A variational approach (à la Yang & Yang) should yield sharp bound on the range of *p* for which BE with real roots only is well-posed.

The range $0 < \alpha < 1$ seems more difficult to study.

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• Transform the logarithmic BE into a Free-Boundary Nonlinear Integral Equation (known as Destri-De Vega).

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• Linearise in the large p limit and do perturbation analysis.

The strategy is standard, the analysis is completely new.

Given $\lambda \vdash N$, let $H = \#\{$ holes greater than the lowest root $\}$. The unknown is a tuple $(\omega, h_1, \dots, h_H, z)$

- $\omega > 0$, the left endpoint of the integration interval $[\omega, +\infty[;$
- $h_1 < \cdots < h_H$ are the holes greater than the lowest root;
- $z : C^1([\omega, \infty[), \text{ strictly monotone, } z(E) \sim E^{\frac{1+\alpha}{2\alpha}}, x \to +\infty.$

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Destri-De Vega Integral Equation II

The Destri-De Vega (DDV) is a free-boundary nonlinear integral equation:

$$1. z(E) = -2p + \int_{\omega}^{\infty} K_{\alpha}(E/y) \left[z(y) - \frac{1}{2} \right] \frac{dy}{y} + H F_{\alpha}\left(\frac{E}{\omega}\right)$$
$$- \sum_{j=1}^{H} F_{\alpha}\left(\frac{E}{h_{k}}\right), \quad K_{\alpha}(x) = x F_{\alpha}'(x) = \frac{\sin\left(\frac{2\pi}{1+\alpha}\right)}{\pi} \frac{x}{1+x^{2}-2x\cos\left(\frac{2\pi}{1+\alpha}\right)}$$
$$2. \left[z(\omega) - \frac{1}{2} \right] = -H$$
$$3. z(h_{j}) = \sigma(j) + \frac{1}{2}, \quad j = 1 \dots H, \quad \sigma(j) = \text{quantum number of } h_{j}$$

Remark. If z is a strictly increasing real analytic function

$$\lim_{\varepsilon \to 0^+} \frac{1}{\pi} \operatorname{Im} \log \left(1 + e^{2\pi i \, z(x+i\varepsilon)} \right) = z - \left[z - \frac{1}{2} \right]$$

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Large p Linearisation = WKB

$$z_{\omega,p}(E) = -2p + \int_{\omega}^{\infty} K_{\alpha}(E/y) z_{\omega,p}(y) \frac{dy}{y}, \ z_{\omega,p}(E) \sim E^{\frac{\alpha+1}{2\alpha}}, x \to \infty.$$

It is a Wiener-Hopf equation, solutions can be expressed via

$$\tau(\xi) = \frac{1}{2\pi i} \int_{\delta - i\infty}^{\delta + i\infty} \frac{\alpha^{\frac{\alpha s}{1 + \alpha}}}{2\sqrt{\pi}(1 + \alpha)^{s - 1}} \frac{\Gamma\left(-\frac{1}{2} - \frac{\alpha s}{1 + \alpha}\right)\Gamma\left(1 - \frac{s}{1 + \alpha}\right)}{s^2 \Gamma(-s)} \xi^{-s} ds, \quad \xi = x/\omega.$$

We discovered a (much more useful) formula in terms of a WKB integral

$$\tau(\xi) = \frac{1}{\pi} \int_{u_{-}}^{u_{+}} \sqrt{u^{2}\xi - u^{2\alpha+2} - 1} \frac{du}{u}, \ \sqrt{\cdots}_{|u=u_{\pm}} = 0.$$

This is a first hint of the ODE/IM correspondence.

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Perturbation/Analytical challenges

We need to analyse integrals like

$$A_{p}[f,\varepsilon] = \int_{1}^{\infty} K_{\alpha}\left(\frac{x}{y}\right) \langle pf(y) + \varepsilon(y) \rangle \frac{dy}{y}, \ \langle z \rangle = z - \left[z - \frac{1}{2}\right]$$
$$B_{p}[f,\varepsilon] = \int_{1}^{\infty} K_{\alpha}\left(\frac{x}{y}\right) \left[pf(y) + \varepsilon(y) - \frac{1}{2}\right] \frac{dy}{y}$$

As an example, we showed that if $f \sim x^{\frac{\alpha+1}{2\alpha}}$ and $\varepsilon, \tilde{\varepsilon}$ are bounded (+ some further hypotheses), then

$$\|B_{\rho}[f,\varepsilon] - B_{\rho}[f,\tilde{\varepsilon}]\|_{\infty} - \frac{\alpha+1}{2\alpha}\|\varepsilon - \tilde{\varepsilon}\|_{\infty} \leq_{f} \frac{\|\varepsilon - \tilde{\varepsilon}\|_{\infty}}{p}$$

 \implies contractiveness of the perturbation operator $B_p[I, \cdot]$ when p is large.

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$$\begin{aligned} A_{\rho}[f,\varepsilon] &= \int_{1}^{\infty} \mathcal{K}_{\alpha}\left(\frac{x}{y}\right) \left\langle \rho f(y) + \varepsilon(y) \right\rangle \frac{dy}{y}, \ \left\langle z \right\rangle = z - \left[z - \frac{1}{2}\right] \\ B_{\rho}[f,\varepsilon] &= \int_{1}^{\infty} \mathcal{K}_{\alpha}\left(\frac{x}{y}\right) \left[\rho f(y) + \varepsilon(y) - \frac{1}{2}\right] \frac{dy}{y} \end{aligned}$$

As an example, we showed that if $f \sim x^{\frac{\alpha+1}{2\alpha}}$ and $\varepsilon, \tilde{\varepsilon}$ are bounded (+ some further hypotheses), then

$$\left| \|B_{p}[f,\varepsilon] - B_{p}[f,\tilde{\varepsilon}]\|_{\infty} - \frac{\alpha+1}{2\alpha} \|\varepsilon - \tilde{\varepsilon}\|_{\infty} \right| \lesssim_{f} \frac{\|\varepsilon - \tilde{\varepsilon}\|_{\infty}}{p}$$

 \implies contractiveness of the perturbation operator $B_p[I, \cdot]$ when p is large.

A family of anharmonic oscillators. Dorey Tateo, BLZ '98

$$-\Psi''(x)+\left(x^{2\alpha}+\frac{\ell(\ell+1)}{x^2}-E\right)\Psi(x)=0, \alpha>1, \ell\geq 0, E\in\mathbb{C}.$$

E is said an eigenvalue if $\exists \Psi \neq 0$ such that

$$\lim_{x\to 0^+}\Psi(x)=\lim_{x\to +\infty}\Psi(x)=0.$$

The spectrum is discrete, simple and positive, ${\it E}_n(\ell), n\in\mathbb{N}$:

$$E_n(\ell) = \left(\frac{2\Gamma(\frac{2\alpha+1}{2\alpha})}{\sqrt{\pi}\Gamma(\frac{3\alpha+1}{2\alpha})}\right)^{-\frac{2\alpha}{\alpha+1}} (4n+2\ell+3)^{\frac{2\alpha}{\alpha+1}} (1+O(n^{-1}))$$

Spectral determinant $D_{\ell}(E)$ is an entire function of order $\frac{1+\alpha}{2\alpha}$

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Monster potentials, BLZ (2003)

1. Let *R* be a monic polynomial of degree *N*. The spectral determinant $D_{\ell}^{R}(E)$ for the potential

$$V^{R} = x^{2\alpha} + \frac{\ell(\ell+1)}{x^{2}} - 2\frac{d^{2}}{dx^{2}} \log R(x^{2\alpha+2})$$

satisfies the BE if the monodromy about the additional poles is trivial for every E.

2. Assuming that the roots of R are distinct, the trivial monodromy is equivalent to the BLZ system

$$\sum_{j \neq k} \frac{z_k \left(z_k^2 + (3+\alpha)(1+2\alpha) z_k z_j + \alpha(1+2\alpha) z_j^2 \right)}{(z_k - z_j)^3} - \frac{\alpha z_k}{4(1+\alpha)} + \Delta(\ell, \alpha) = 0 , \quad k = 1, \dots, N.$$

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Rational extensions of the harmonic oscillator

• A rational extension of degree N is a potential

$$V^{U}(t) = t^{2} - 2 \frac{d^{2}}{dt^{2}} \ln U(t),$$

where U a polynomial of degree N such that all monodromies of $\psi''(t) = (V^U(t) - E)\psi$ are trivial for every E.

• Oblomkov's theorem (1999)

$$U \propto U^{\lambda} := Wr[H_{\lambda_1+j-1}, \dots, H_{\lambda_j}], \text{ for a } \lambda := (\lambda_1, \dots, \lambda_j) \vdash N.$$

Large momentum limit of Monster Potentials

(2) (Conditional) Theorem, M. - Conti 2021/2022

• We noticed that in the large momentum multi-scale limit, monster potentials converge to rational extensions of the harmonic oscillator:

Assume there exists a sequence R_ℓ of monster potentials with $\ell\to\infty,$ then – up to subsequences –

$$z_k = \frac{\ell^2}{\alpha} + \frac{(2\alpha+2)^{\frac{3}{4}}}{\alpha} v_k^{\lambda} \ell^{\frac{3}{2}} + O(\ell), \ k = 1, \dots, N$$

where v_k^{λ} are the roots of U^{λ} .

• (If a monster potential with a such an asymptotics exists and) $D_{\ell}^{\lambda}(E)$ is the corresponding spectral determinant, then

$$D^{\lambda}(E;\ell) = Q^{\lambda}(E/\eta;p), \ P = \frac{2\ell+1}{\alpha+1} \text{ and } \eta = \left(\frac{2\sqrt{\pi}\,\Gamma\left(\frac{3}{2}+\frac{1}{2\alpha}\right)}{\Gamma\left(1+\frac{1}{2\alpha}\right)}\right)$$

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Let $\lambda \vdash N$, assume U^{λ} has N distinct zeroes (see conjecture by Felder-Hemery-Veselov 2010). Consider the Jacobian

$$J_{ij}^{\lambda} = \delta_{ij} \left(1 + \sum_{l
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The Big ODE/IM Conjecture

If a QFT is Bethe Integrable then the corresponding solutions of the Bethe Equations are **spectral determinants of linear differential operators**.

 \rightarrow Bethe Roots are eigenvalues of a (possibly self-adjoint) differential operator (cf. Hilbert-Pólya Conjecture).

M - Raimondo (- Valeri) ('16,'17, '20, ongoing) after Feigin-Frenkel (2011)

 $\widehat{\mathfrak{g}}$ an affine Kac-Moody Lie-algebra and ${}^L\widehat{\mathfrak{g}}$ the Langlands dual,

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The ODE/IM correspondence for Quantum KdV is just a tiny piece of an enormous field of research of which we know a lot but still very little.

- How do we guess which ODE (if any) corresponds to a given Quantum Field Theory?
- Once, they are found, how do we prove them?
- Why the ODE/IM correspondence? Can we find a theory? Why is the nonlinear Stokes phenomenon that important?

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